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## 18.6-year lunar nodal tides from altimeter data

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## ABSTRACT

Harmonic analyses of 16-year long time series of collocated TOPEX, Poseidon and Jason-1 altimeter data were carried out in the Pacific and western Atlantic Oceans and their marginal seas. These time series are sufficiently long to adequately separate the 18.6-year nodal satellites  $Q_{1n}$ ,  $O_{1n}$ ,  $K_{1n}$ ,  $M_{2n}$  and  $K_{2n}$  from their parent constituents  $Q_1$ ,  $O_1$ ,  $K_1$ ,  $M_2$  and  $K_2$ . Editing criteria were used to eliminate results in areas where these satellites are weak (i.e., smaller than their formal error estimates), or where they are strongly affected by aliased low-frequency signals (e.g. in the Kuroshio, in the Gulfstream and in their extension regions). As expected from tidal theory, the phases of the altimetry-derived nodal satellites agree reasonably well with the phases of their parents. However, due to their relatively small amplitudes and the remaining influence from low-frequency aliased signals, the altimeter observed amplitude ratios between the nodal satellites and their parent constituents tend to exceed the values predicted by the theory.

Examination of diurnal and semidiurnal nodal amplitudes in select coastal areas and marginal seas around the Pacific and the western Atlantic Ocean allowed the assignment of a *nodal character* to regions, which were each classified as *nodal diurnal*, *nodal semidiurnal*, or *nodal mixed*, based on the nodal amplitudes in each band. While the areas with predominant diurnal tides are all *nodal diurnal*, the small nodal ratio of 0.037 for  $M_{2n}$  resulted in some regions with strong  $M_2$  tides being classified as *nodal diurnal* or *nodal mixed*. The amplitude ratio between  $K_{2n}$  and  $K_2$  is 0.30, making the  $K_{2n}$  amplitudes sometimes comparable to those of  $M_{2n}$ . However, this effect was not sufficient to make all the areas with dominant  $M_2$  to be dominant *nodal semidiurnal*.

The observed amplitudes of the 18.6-year nodal constituent  $M_n$  are relatively small, 1.5–3.5 cm. These values significantly exceed its theoretical amplitudes, which are less than 1 cm almost everywhere. The analysed signals at  $M_n$  frequency are therefore of mostly non-tidal origin, part of the broad-band decadal ocean variability.

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## 1. Introduction

The launch of the TOPEX/Poseidon satellite in 1992 with accurate radar altimeters has provided a new and very important data set for tidal analysis of global sea levels (e.g., Schrama and Ray, 1994; Ray, 1998; Cherniawsky et al., 2001) and resulted in very significant improvements in the accuracy of tidal prediction and of numerical tidal models (e.g., Egbert et al., 1994; Le Provost et al., 1998; Ray, 1999; Egbert and Erofeeva, 2002; Foreman et al., 2000, 2006). In 2002, the TOPEX/Poseidon satellite was replaced in the same orbit by Jason-1, also with a very accurate altimeter, thus extending the along-track time series of sea level observations to more than 16 years. The longer time series not only

increases the accuracy of computed constituents but also allows analyses of minor constituents with frequencies relatively close to the major ones.

Modern harmonic analyses of tidal heights (Godin, 1972; Foreman, 1977, hereafter F77) follow the development of tidal potential theory by Doodson (1921) and express each constituent frequency as a linear superposition of six astronomical forcing harmonics,  $\omega_k = l_1\tau + l_2s + l_3h + l_4p + l_5n' + l_6p'$ , where  $l_i$  are integers and  $\tau$ ,  $s$ ,  $h$ ,  $p$ ,  $-n'$  and  $p'$  are the mean rates of change of lunar time (with a mean period of 24.84 h) and of the longitudes of the moon (27.3 days), the sun (365.24 days), the lunar perigee (8.85 years), the moon's ascending node (18.61 years) and the solar perigee (20 392 years), respectively.

Usually when time series are limited to a few years, only the frequencies that differ by at least one multiple of the third harmonic ( $h$ ) can be included in the analyses, provided that their expected amplitudes exceed background noise at the same frequency. But for longer time series, with durations approaching

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or exceeding the 18.61-year nodal period, it is also possible to calculate the amplitudes and phases of nodal satellites of major constituents, which were previously included as part of the nodal corrections (Godin, 1972, 1986). The 18.6-year nodal variations have a direct effect on annual means of high water, low water and tidal range at the coast (e.g., Kaye and Stuckey, 1973; Amin, 1993), thus requiring detailed analyses of long-term sea level records.

One of the earliest analyses of nodal constituents was provided by Doodson (1924), who examined 19-year long variations in amplitudes of  $M_2$  and  $N_2$  in tidal records from Saint John, New Brunswick, and Bombay, India. His results for  $M_2$  at Saint John showed a reasonably regular 19-year cycle, as expected from the nodal modulation. However variations in  $N_2$  showed a more complex behaviour, likely due to the 9-year cycle in the longitude of lunar perigee. More recently, Amin (1985) analysed an 18.61-year long time series of sea level at stations on the west coast of Great Britain. His analyses of the 1963–1981 time series at Liverpool and Newlyn showed good agreement with theoretical nodal amplitude ratios for the major constituents, though the 0.032 ratio obtained for  $M_2$  at Liverpool was smaller than the theoretical value of 0.037.

Another example of analyses of two 19-year long time series (1939–1957, 1958–1976) of hourly sea level in Victoria Harbour was described by Foreman and Neufeld (1991), who included more than 500 astronomical and shallow water constituents, of which about 180 (including 20 low-frequency and 47 shallow-water) constituents were found to be strong enough to be listed in the paper. In general, the observed nodal satellite ratios at Victoria agreed well with the tidal equilibrium values, though the ratios for nodal satellites of  $N_2$  and  $M_2$  were somewhat smaller, 0.028 and 0.034, than the 0.037 value predicted by the theory. These results demonstrate that, provided a tidal constituent amplitude lies above the noise level, its relative amplitude and phase generally agree with tidal potential theory predictions, thus confirming the usefulness of satellite correction algorithms employed by Godin (1972, 1986) and F77. However, local effects, such as friction in shallow water, or tidal resonance, also affect the observed nodal ratios, thus making them different from their astronomical values (e.g., Ku et al., 1985).

The 18.6-year constituent  $M_n$  was also included in the Foreman and Neufeld (1991) analyses and appears to be nonstationary. Its amplitudes for the two 19-year periods were 1.2 and 1.4 cm, respectively, which are about twice as high as the equilibrium value of 0.6 cm for this latitude (e.g., Pugh, 1987). The computed  $M_n$  phase values were also different for the two periods,  $269^\circ$  and  $200^\circ$ , respectively, i.e.  $91^\circ$  and  $160^\circ$  ahead of the  $M_n$  equilibrium phase. However, when using all data from the combined period of 1939–1976, the analysed  $M_n$  amplitude is somewhat smaller, 1.1 cm, and its phase of  $300^\circ$  is closer to the equilibrium phase. It is then natural to assume that for sufficiently long sea level timeseries, the analysed  $M_n$  amplitude and phase may be close to their equilibrium values. Such a result was indeed presented by Trupin and Wahr (1990), who used yearly averages of global tide gauge measurements from 260 stations with at least 19 years of data to show that the 18.6-year peak in the amplitude spectrum of “stacked” data (see Trupin and Wahr, 1990 for an explanation of “stacking”) fits well the equilibrium spectral peak for  $M_n$ .

The 18.6-year modulation of diurnal and semidiurnal tides is of special interest because in many parts of the ocean it exerts a strong effect on tidal currents and vertical mixing. For example, Loder and Garrett (1978) and Garrett (1979) considered a simple model of vertical mixing on the shelf and found that even small, of the order of  $\varepsilon$ , variation in tidal currents can have a very noticeable effect on sea surface temperature (SST),  $\delta T \sim (4\varepsilon/3)\Delta T$ , where  $\Delta T$  is the top to bottom temperature difference. The Loder and Garrett (1978) analyses of coastal SSTs on the east and west

coasts of North America yielded amplitudes of the 18.6-year signal between 0.16 and 0.64 °C and phases in agreement with their model predictions.

Royer (1993) analysed ocean and air temperatures in the Gulf of Alaska and found that up to 30% of their low-frequency variations can be accounted for with the 18.6-year nodal signal. The long-term variations in tidal mixing and ocean temperatures can have important consequences for biological productivity and fisheries (e.g., Parker et al., 1995). Munk et al. (2002) discuss the role of such variations in tidal mixing, which include nodal modulations, on the oceanic poleward heat transport and climate. Ray (2007) extended the Loder and Garrett (1978) calculations of coastal SSTs and also provided a detailed analysis of the various hypotheses and mechanisms linking the nodal tides to decadal ocean variability. He concluded that the most likely process responsible for this variability is the 18.6-year modulation of diurnal tidal currents and vertical mixing.

Osafune and Yasuda (2006) found the bi-decadal oscillations in the northwestern Pacific Ocean to be synchronized with the 18.6-year nodal cycle, while Yasuda et al. (2006) linked this cycle to Pacific Ocean climate variability and to the Pacific Decadal Oscillation (PDO). Hasumi et al. (2008) used a coupled ocean-atmosphere climate model to show that the 18.6-year cycle in tidal mixing around the Kuril Islands leads to coastal trapped waves propagating along the perimeter of the ocean, leading to 18.6-year periodicity in the Pacific Ocean variability. McKinnell and Crawford (2007) have found the nodal variations in northeast Pacific Ocean temperatures to be linked with the Pacific–North America (PNA) index, with the latter lagging the nodal signal by about 2 years. Nodal modulations also affect high latitude oceans temperatures and salinities (Yndestad et al., 2008) and the Arctic Ocean climate (Yndestad, 2006).

The effects of nodal tides on mixing can be readily investigated in numerical tidal models. For example, Foreman et al. (2006) found the nodal modulation to cause basin-wide variations of approximately 19% in the Bering Sea tidal energy flux, with even stronger variations in the straits of the Aleutian Islands, such as Seguam Pass, and around capes. It is therefore of interest to identify other areas where nodal modulation of diurnal or semidiurnal tides can have significant effects on ocean mixing and circulation, while examining the accepted wisdom that such effects are mainly due to diurnal tides.

In this paper, we carry out harmonic analyses of more than 16-year long records of TOPEX/Poseidon and Jason-1 (TPJ) along-track altimetry data in the Pacific and the western Atlantic Oceans (20S to 66N, 100E to 40W) in order to compute the nodal satellites of five major tidal constituents  $Q_1$ ,  $O_1$ ,  $K_1$ ,  $M_2$  and  $K_2$ . While these analyses are performed over the complete area, our attention is further focused on fifteen select regions, which include certain coastal areas and marginal seas. Some of these regions exhibit strong diurnal tides, while others show semi-diurnal or mixed tides. In the next two sections, we describe in some detail the processing and harmonic analysis procedures, with specific attention given to possible aliasing due the 9.916-day TPJ repeat period. The analyses results and spatial variations in nodal amplitudes and phases are presented and mean nodal ratios from altimetry are compared to the theoretical values. This comparison is carried out for the complete area and for each of the 15 regions.

## 2. Selection and preprocessing of altimeter data

Along-track altimetry data were extracted for the area of study using the Institute of Ocean Sciences (Sidney, BC) mirror of the Radar Altimetry Data System (RADS: <http://rads.tudelft.nl/rads/rads.shtml>) and associated software (Schrama et al., 2000; Naeije

et al., 2002; Scharroo, 2008). These data include 364 cycles (September 23, 1992–August 11, 2002) from TOPEX and Poseidon along their original ground tracks and 260 cycles (January 16, 2002–February 5, 2009) from Jason-1 altimeter, along the same ground tracks, thus giving a total of 624 cycles in the combined (TPJ) data set. However, during the tandem phase in 2002, the TOPEX/Poseidon satellite was in the same orbit and just 71 s behind Jason-1, so that the number of independent cycles is 603.

The default RADS geophysical corrections and flag selection (Scharroo, 2008) were used for each altimeter, except we retained the ocean tides in the data and used load tides derived from a current version of the GOT tidal model (GOT4.7), described in Ray (1999). We also ignored the somewhat conservative “radiometer land flag”, thus allowing inclusion of data closer to land, while relying on other flags and editing to screen questionable data. RADS default parameters include reference frame offsets, ensuring that the three data sets are related at each location to the same reference level. Appendix A describes the editing and collocation of along-track data for each altimeter. Collocated data from each of the three altimeters were spliced together in chronological order, forming TPJ time series that are about 16.3 years long (September 23, 1992–February 5, 2009).

### 3. Harmonic analysis and aliasing

We calculate select tidal constants at each along-track TPJ location using the well known harmonic analysis procedure (F77), subsequently modified for TOPEX/Poseidon data (Cherniawsky et al., 2001, hereafter CFCH). As the time period and the number of observations (cycles) increase, it becomes possible to include in the analysis a larger number of constituents and satellites.

After several tests, we selected 36 constituents (including  $Z_0$ ), based on relative amplitudes from tidal potential theory (Doodson, 1921; Godin, 1972; Pugh, 1987; F77; Foreman and Neufeld, 1991) and a posteriori plots of amplitudes in the area of study. These constituents are listed in Table 1 with their six Doodson (1921) numbers, frequencies, and the main (lowest-frequency) periods  $T_{al}$  of aliased signals due to the TPJ repeat period of 9.916 days (Parke et al., 1987; Ray, 1998, CFCH). The more than 16-year ( $\sim 5970$  days) long time series are long enough to allow good separation of potentially aliased pairs, such as  $S_{sa}-K_1$  and  $P_1-K_2$ , but still somewhat too short to fully resolve the 18.6-year nodal satellites, denoted in Table 1 with a subscript n.

Examination of along-track error covariance matrices (Appendix B) shows only minor contributions from off-diagonal elements to constituent amplitude errors, even for the 18.6-year satellites. Mean values of error estimates  $\sigma_a$  ( $=\sigma_k$  in Appendix B) are also listed in Table 1, together with the number of remaining valid locations  $N_v$  after the editing (explained below). However, these estimates do not (and cannot) include contributions from low-frequency aliased signals that are not in the analysis (e.g., Munk and Cartwright, 1966; Ray, 1998). For the frequency range in question, the oceanic sea level power spectrum is red. Therefore, analysed constituents with longer aliased periods  $T_{al}$  (Table 1) are expected to have larger errors due to such aliasing. In particular,  $K_1$  ( $T_{al}=173.3$  days) and  $K_{1n}$  (177.9 days) tend to show short-scale along-track amplitude variations that exceed their formal error estimates, which may be, at least in part, due to such low-frequency aliasing.

To minimize the effects of aliasing and of strong non-tidal signals, we adopted the following edit criteria when plotting, or tabulating computed amplitudes  $A$  and Greenwich phase lags  $G$  of the analysed constituents.

**Table 1**

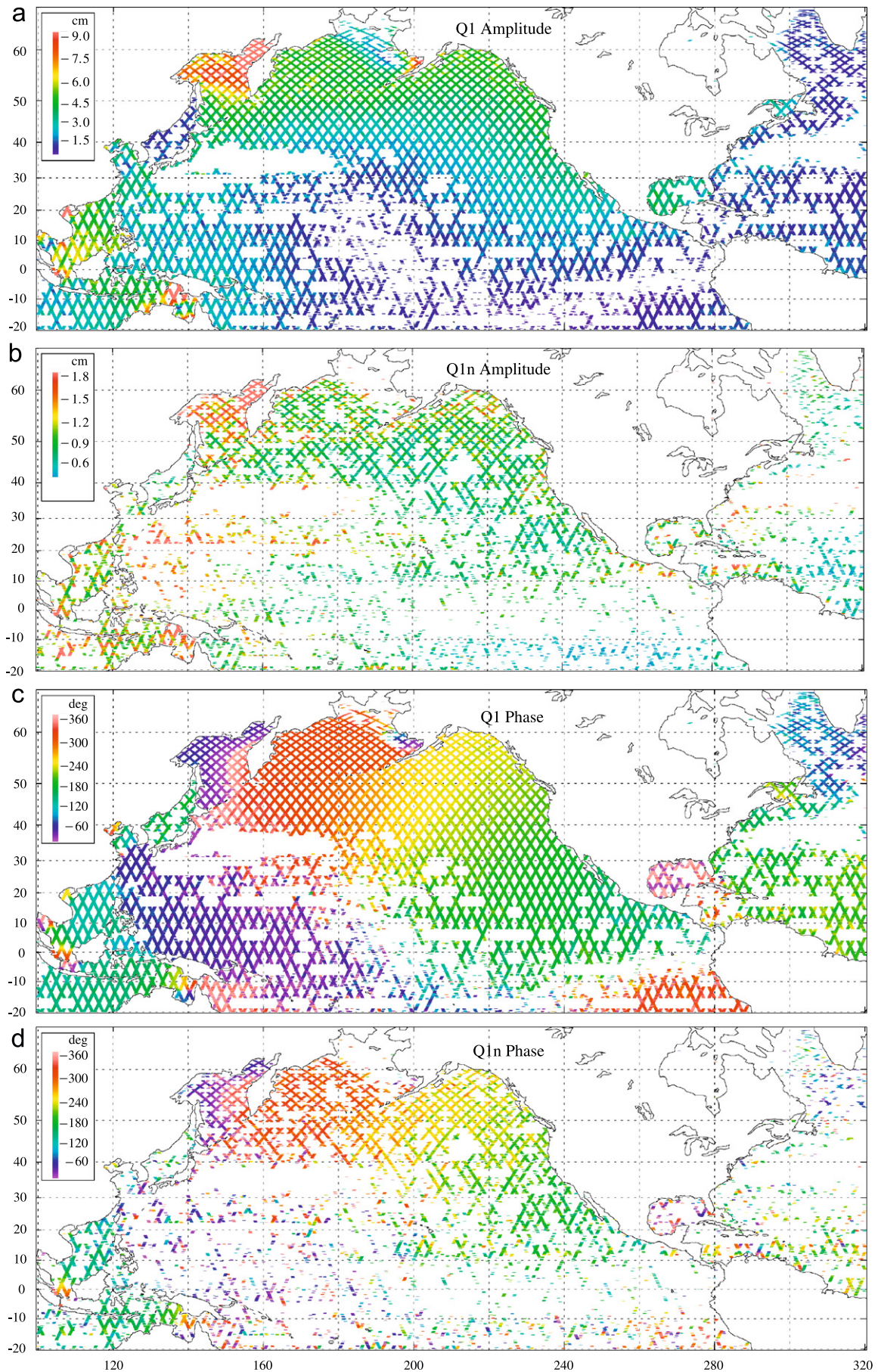
Tidal constituents selected for harmonic analysis, listed in order of frequency, with Doodson numbers and main (lowest-frequency) alias periods  $T_{al}$ .

	Doodson numbers	Frequency (h <sup>-1</sup> )	$T_{al}$ (day)	$\bar{A}_a$ (cm)	$\bar{\sigma}_a$ (cm)	$N_v$
$Z_0$	0 0 0 0 0 0	0.000000000	9.916	4.0	0.4	180350
$M_n$	0 0 0 0 1 0	0.000006129	9.930	<b>2.5</b>	0.8	138964
$S_a$	0 0 1 0 0-1	0.000114074	10.192	4.9	0.8	181567
$S_{sa}$	0 0 2 0 0 0	0.000228159	10.485	1.9	0.8	135022
$M_n$	0 1 0-1 0 0	0.001512152	15.490	0.9	0.6	73028
$M_{sf}$	0 2-2 0 0 0	0.002821933	30.189	0.9	0.7	11623
$M_f$	0 2 0 0 0 0	0.003050092	36.168	1.3	0.7	125292
$\sigma_1$	1-3 2 0 0 0	0.035908722	21.812	0.9	0.7	26024
$Q_{1n}$	1-2 0 1-1 0	0.037212374	68.682	<b>1.1</b>	0.8	39248
$Q_1$	1-2 0 1 0 0	0.037218503	69.383	<b>2.8</b>	0.9	124510
$\rho_1$	1-2 2-1 0 0	0.037420874	104.648	1.1	0.7	47120
$O_{1n}$	1-1 0 0-1 0	0.038724526	46.015	<b>2.7</b>	0.8	135078
$O_1$	1-1 0 0 0 0	0.038730654	45.706	<b>11.5</b>	0.8	185521
$NO_1$	1 0 0 1 0 0	0.040268594	23.775	1.3	0.6	84593
$\pi_1$	1 1-3 0 0 1	0.041438513	71.514	1.0	0.7	21505
$P_1$	1 1-2 0 0 0	0.041552587	88.925	6.1	0.9	153371
$S_1$	1 1-1 0 0 1	0.041666672	117.545	1.8	1.1	36958
$K_1$	1 1 0 0 0 0	0.041780746	173.322	<b>17.8</b>	0.9	177027
$K_{1n}$	1 1 0 0 1 0	0.041786875	177.856	<b>3.3</b>	0.8	97998
$J_1$	1 2 0-1 0 0	0.043292898	32.763	1.4	0.7	90542
$2N_2$	2-2 0 2 0 0	0.077487097	22.534	1.4	0.7	114749
$\mu_2$	2-2 2 0 0 0	0.077689468	20.311	1.6	0.7	124293
$N_2$	2-1 0 1 0 0	0.078999249	49.548	8.2	0.8	186348
$\nu_2$	2-1 2-1 0 0	0.079201620	65.251	2.0	0.8	123636
$M_{2n}$	2 0 0 0-1 0	0.080505272	62.648	<b>1.9</b>	0.8	114094
$M_2$	2 0 0 0 0 0	0.080511401	62.076	<b>37.9</b>	0.9	190562
$\lambda_2$	2 1-2 1 0 0	0.081821182	21.033	1.0	0.7	18709
$L_2$	2 1 0-1 0 0	0.082023553	20.640	1.4	0.7	101384
$T_2$	2-2 3 0 0 1	0.083219259	50.626	1.3	0.7	94039
$S_2$	2 2-2 0 0 0	0.083333333	58.772	13.8	0.9	185009
$K_2$	2 2 0 0 0 0	0.083561492	86.661	<b>4.2</b>	0.9	165825
$K_{2n}$	2 2 0 0 1 0	0.083567624	87.780	<b>1.7</b>	0.8	107801
$\eta_2$	2 3 0-1 0 0	0.085073644	40.400	0.8	0.6	25892
$MO_3$	3-1 0 0 0 0	0.119242055	26.324	1.0	0.8	7682
$M_3$	3 0 0 0 0 0	0.120767101	38.079	0.8	0.6	25407
$M_4$	4 0 0 0 0 0	0.161022801	31.038	1.2	0.8	16376

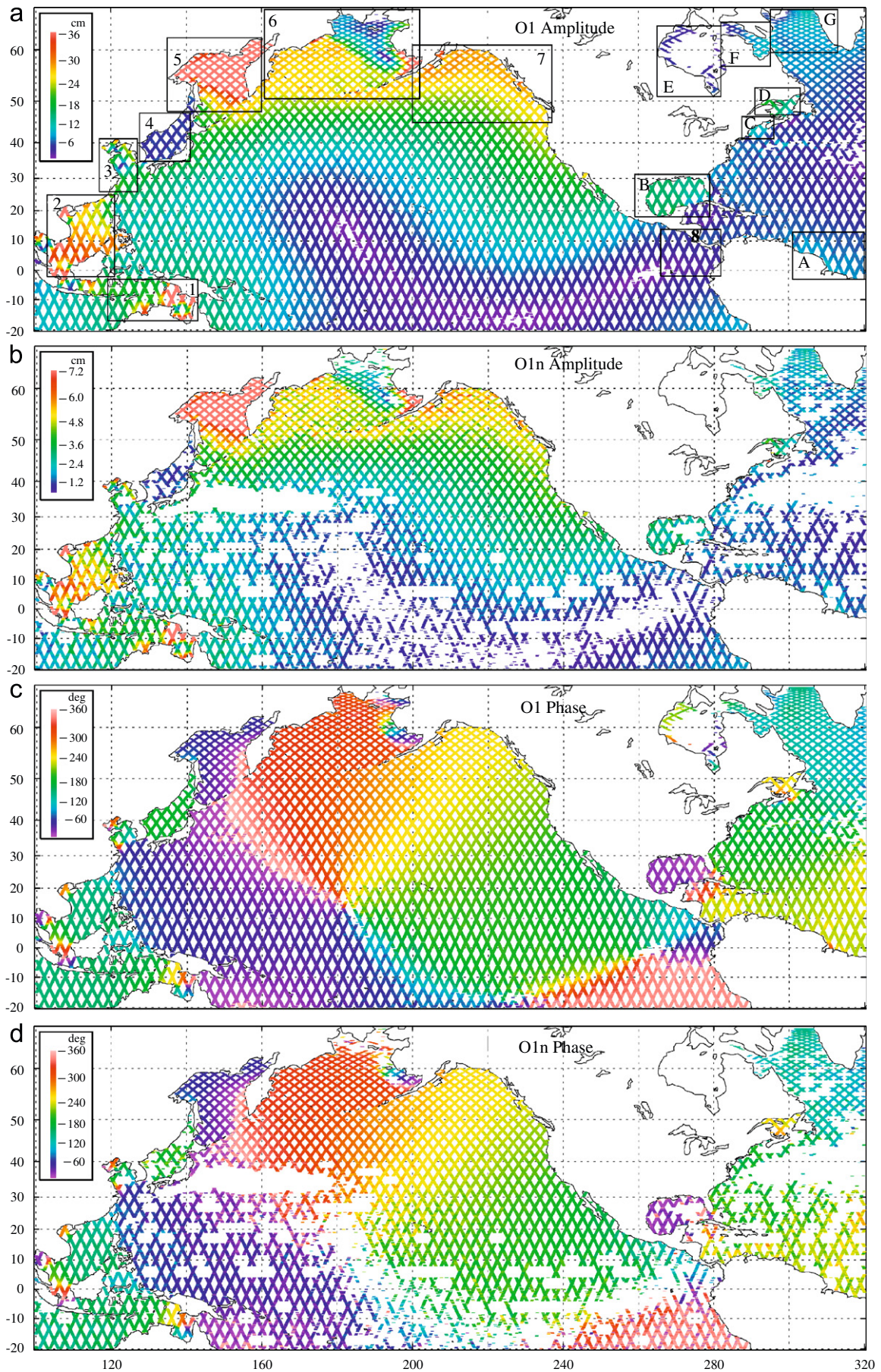
Mean amplitudes  $\bar{A}_a$  from harmonic analysis and their mean error estimates  $\bar{\sigma}_a$  are also shown, computed from data at  $N_v$  valid locations after applying the three editing criteria. Amplitude values marked in bold are for the constituents that are the subject of this work.

1. We require the analysed amplitudes to be larger than their error estimates:  $A > \sigma$ .
2. We require the computed complex amplitudes  $C = A \exp(iG)$  on ascending and descending passes to agree near pass cross-overs. Using  $r = 2|C_{asc} - C_{desc}| / |C_{asc} + C_{desc}|$ , where  $C_{asc}$  and  $C_{desc}$  are average complex amplitudes at two ascending and two descending along-track locations near a crossover (Fig. A1), we exclude constituents near a crossover, and up to half the distance between the crossovers, when  $r > 0.5$ .
3. We exclude locations where the root-mean-square sea level variability of the residual signal after detiding exceeds 1.6 m.

These edit criteria remove most, though not all suspect data. The number of remaining valid data locations ( $N_v$  in Table 1) depends on the relative amplitudes and susceptibility to aliasing of each constituent. For consistency, the last three columns in Table 1 were produced using all three edit criteria. However, we only used the first criterion when plotting amplitudes and phases of the three strongest constituents  $O_1$ ,  $K_1$  and  $M_2$ .



**Fig. 1.** Computed (a,b) amplitudes and (c,d) phases of (a,c)  $Q_1$  and (b,d)  $Q_{1n}$ .



**Fig. 2.** Computed (a,b) amplitudes and (c,d) phases of (a,c)  $O_1$  and (b,d)  $O_{1n}$ .

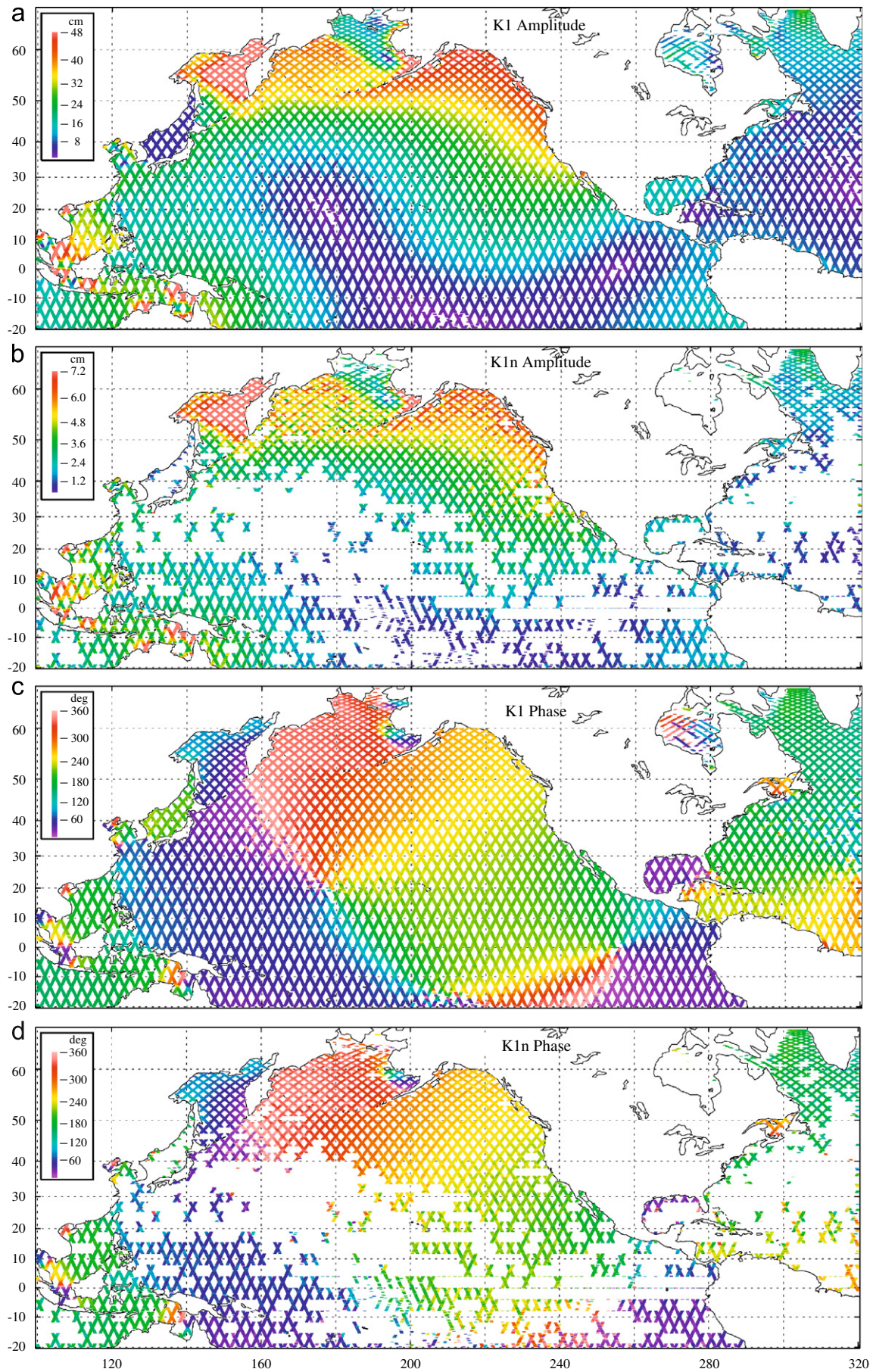
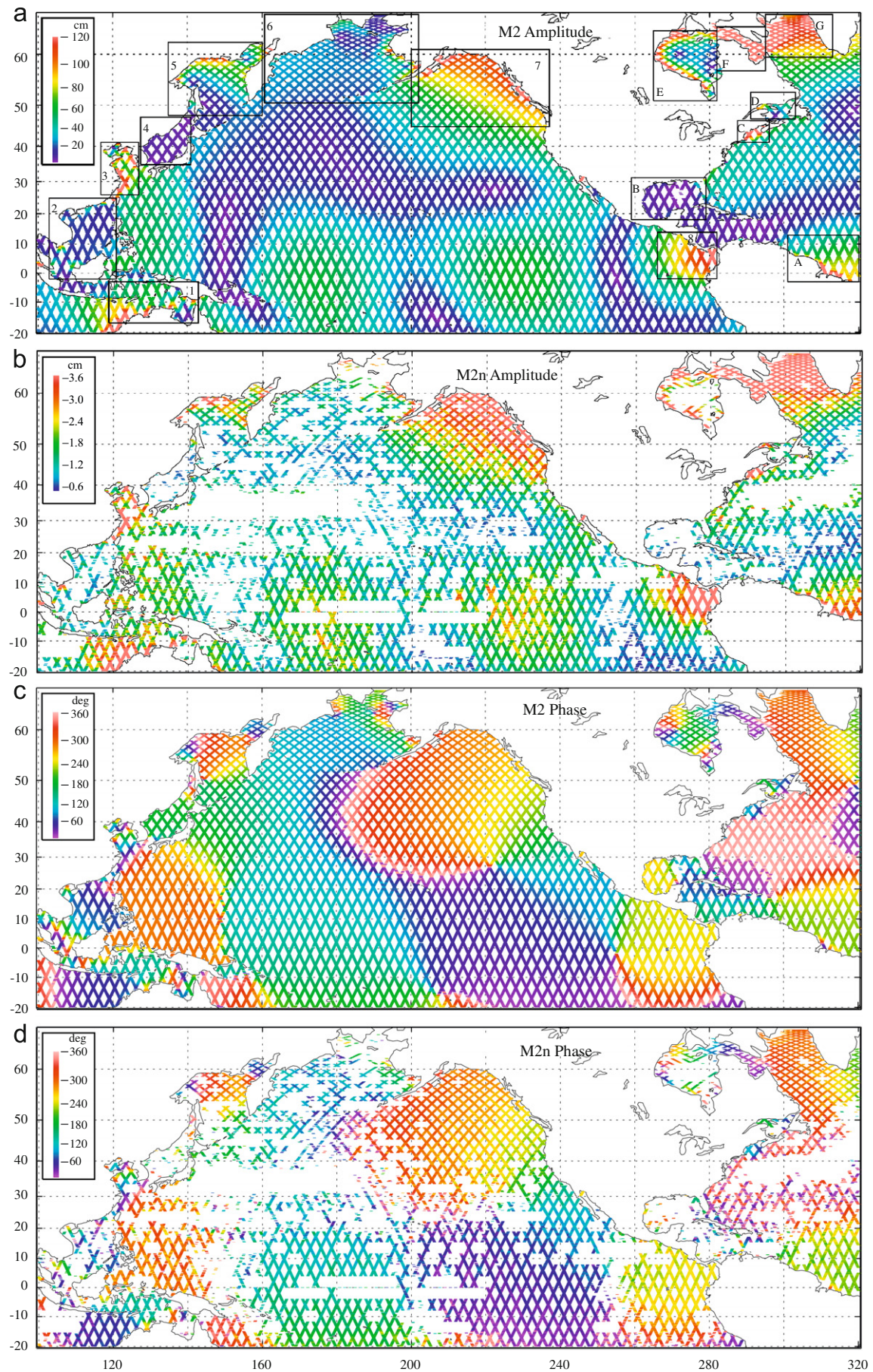


Fig. 3. Computed (a,b) amplitudes and (c,d) phases of (a,c)  $K_1$  and (b,d)  $K_{1n}$ .



**Fig. 4.** Computed (a,b) amplitudes and (c,d) phases of (a,c)  $M_2$  and (b,d)  $M_{2n}$ .

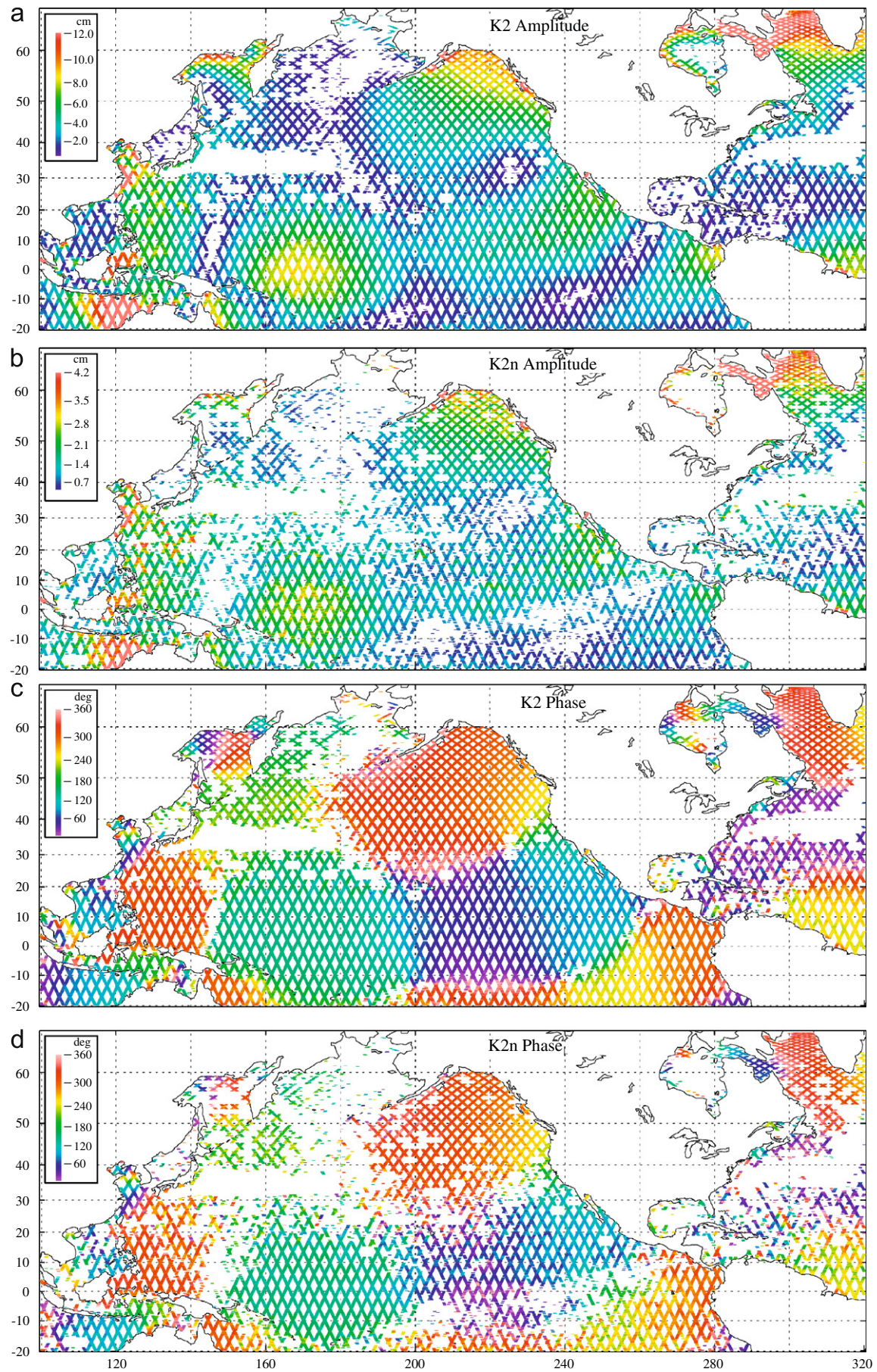


Fig. 5. Computed (a,b) amplitudes and (c,d) phases of (a,c)  $K_2$  and (b,d)  $K_{2n}$ .

**Table 2**

Average amplitude ratios and phase differences (degrees) between nodal satellites and their parent constituents computed in each region (1–8 in the Pacific and A–G in the Atlantic) after applying the three editing criteria and excluding outliers outside two standard deviations.

Region	$Q_{1n}/Q_1$	$O_{1n}/O_1$	$K_{1n}/K_1$	$M_{2n}/M_2$	$K_{2n}/K_2$
Theory	0.1884	0.1885	0.1356	0.0373	0.2980
Full	0.0 0.43 (0.34)	0.0 0.200 (0.038)	0.0 0.147 (0.034)	0.0 0.044 (0.015)	0.0 0.39 (0.20)
1	–1.3 (38.3) 0.260 (0.072) –5.0 (18.9)	0.3 (12.1) 0.190 (0.017) –0.2 (6.3)	0.5 (12.9) 0.138 (0.014) –0.5 (5.9)	–0.6 (17.5) 0.0400 (0.0076) –0.6 (11.0)	1.1 (23.5) 0.308 (0.060) 4.0 (16.5)
2	0.237 (0.067) 4.2 (22.3)	0.185 (0.015) 1.0 (5.4)	0.135 (0.018) 0.1 (7.7)		
3		0.191 (0.035) 1.0 (8.8)	0.139 (0.019) –0.8 (11.9)	0.0376 (0.0062) 1.3 (13.8)	0.330 (0.076) –3.0 (11.2)
4		0.207 (0.043) 5.8 (17.7)			
5	0.222 (0.063) –3.3 (18.6)	0.194 (0.010) –1.0 (5.2)	0.147 (0.012) –0.7 (5.8)	0.043 (0.011) –3.9 (16.3)	0.39 (0.12) 0.3 (22.5)
6	0.225 (0.052) –4.3 (18.2)	0.201 (0.040) –0.5 (7.7)	0.144 (0.021) –0.9 (7.5)	0.048 (0.014) –1.2 (19.0)	
7	0.231 (0.052) 0.9 (15.4)	0.190 (0.010) 1.5 (3.4)	0.139 (0.008) 0.0 (3.2)	0.0376 (0.0041) –0.1 (6.9)	0.314 (0.035) 0.3 (9.4)
8			0.157 (0.043) 5.3 (14.5)	0.0392 (0.0043) –0.3 (5.6)	0.259 (0.055) –1.9 (14.1)
A		0.190 (0.039) –1.7 (10.8)	0.158 (0.033) 0.3 (18.4)	0.0379 (0.0044) –0.9 (7.1)	0.269 (0.061) –6.5 (15.6)
B		0.192 (0.040) –0.6 (13.5)	0.150 (0.051) 3.1 (29.0)		
C		0.183 (0.041) 10.6 (8.8)	0.155 (0.031) –17.8 (14.8)	0.0351 (0.0038) 1.3 (5.9)	0.340 (0.077) 10.7 (20.5)
D		0.188 (0.027) –2.4 (7.2)	0.143 (0.029) 0.6 (9.3)	0.044 (0.012) 0.4 (9.5)	0.39 (0.16) 7.3 (27.2)
E				0.044 (0.012) –2.6 (16.4)	0.40 (0.12) 2.4 (15.6)
F		0.291 (0.050) –2.3 (13.7)	0.198 (0.057) 17.7 (12.9)	0.0363 (0.0028) –1.0 (4.3)	0.305 (0.023) –1.2 (4.8)
G		0.213 (0.035) –1.5 (10.3)	0.144 (0.026) 2.4 (9.2)	0.0402 (0.0039) –2.6 (7.0)	0.312 (0.038) 0.8 (7.6)

Values in parentheses are the corresponding standard deviations.

#### 4. Analysis results

Figs. 1–5 display the amplitudes and phases for five major constituents  $Q_1$ ,  $O_1$ ,  $K_1$ ,  $M_2$  and  $K_2$  and their 18.6-year nodal satellites  $Q_{1n}$ ,  $O_{1n}$ ,  $K_{1n}$ ,  $M_{2n}$  and  $K_{2n}$ . In areas not excluded by the above edit criteria, these figures show remarkable agreement in spatial patterns for each parent–satellite pair. This spatial agreement between the nodal satellites and their parents is as expected from their very close frequencies (separated by only  $1/18.6 \text{ year}^{-1}$ ), which govern coastal and shelf wave response and propagation speeds.

Average amplitude ratios between each of these satellites and their parents are listed in Table 2 for the complete area shown in Figs. 1–5, as well as for 15 select regions, eight of them in Pacific Ocean (numbered 1–8 in Figs. 2 and 4) and seven in the western Atlantic Ocean (A–G). Some of the regions are dominated by diurnal tides, others by semidiurnals. Hence, the analysed amplitude ratios tend to be close to their theoretical values, but not everywhere. In general, they exceed these ratios due to likely contributions from aliased non-tidal signals, and more so for the weaker satellites. Phase values of nodal satellites also agree quite well with their parent constituents, even for the weaker satellites, such as  $M_{2n}$  (Fig. 4), though there are some differences, of the order of  $\pm 10^\circ$ .

Regional mean amplitudes of the four strongest constituents  $O_1$ ,  $K_1$ ,  $M_2$  and  $K_2$  and of their nodal satellites are summarized in Table 3 (see also Figs. 1, 2, 4 and 5). Instead of using the analysed amplitudes of nodal satellites, which are subject to contamination

**Table 3**

Average amplitudes (in cm) of diurnal  $O_1$  and  $K_1$  and semidiurnal  $M_2$  and  $K_2$  constituents in each region and their nodal satellite amplitudes from theoretical ratios (Table 2).

Region	$O_1$	$O_{1n}$	$K_1$	$K_{1n}$	$M_2$	$M_{2n}$	$K_2$	$K_{2n}$	
1	25.1	4.7	35.3	4.8	65.0	2.4	8.7	2.6	D
2	28.5	5.4	33.9	4.6	22.8	0.9	3.1	0.9	D
3	16.5	3.1	22.3	3.0	86.9	3.2	9.5	2.8	M
4	5.1	1.0	5.4	0.7	5.8	0.2	1.1	0.3	D
5	38.4	7.2	50.4	6.8	38.4	1.4	4.6	1.4	D
6	21.5	4.1	32.3	4.4	28.3	1.1	2.2	0.7	D
7	25.0	4.7	40.0	5.4	82.4	3.1	7.3	2.2	D
8	2.8	0.5	9.6	1.3	87.7	3.3	5.8	1.7	S
A	7.1	1.3	7.5	1.0	63.9	2.4	5.7	1.7	S
B	13.5	2.5	14.5	2.0	7.5	0.3	1.5	0.4	D
C	9.3	1.8	11.5	1.6	102.9	3.8	4.3	1.3	S
D	14.3	2.7	14.1	1.9	38.6	1.4	4.0	1.2	D
E	3.6	0.7	12.1	1.6	61.7	2.3	7.6	2.3	S
F	8.0	1.5	16.4	2.2	274.4	10.2	26.6	7.9	S
G	8.9	1.7	17.3	2.3	114.0	4.3	11.4	3.4	S

The corresponding nodal character of each region is marked in the last column, “D” for diurnal-dominant nodal variations, “S” for semidiurnal, or “M” for mixed.

by noise and reduced number of observations ( $N_v$ ), we computed their amplitudes using parent constituent amplitudes multiplied by the theoretical ratios (Table 2). The source of nodal variations in each region depends on whether diurnal or

semidiurnal tides are prevalent and on the nodal ratio of each constituent. In analogy with major constituents, we define a *nodal character* for each region, marked in Table 3 as “D” for diurnal-dominant nodal variations, “S” for semidiurnal and “M” for mixed nodal tides.

Six out of the eight regions in the Pacific Ocean are classified as *nodal diurnal*, with only the Gulf of Panama (region 8) exhibiting a *nodal semidiurnal* character. While  $M_2$  is the dominant constituent in the Yellow and East China Seas (region 3), the relatively small nodal ratio (0.037) for  $M_{2n}$  and the fairly significant diurnal tides in this area cause this region to be classified as *nodal mixed*. Both the Sea of Okhotsk (region 5) and the Bering Sea (region 6) are dominated by diurnal tides and are therefore *nodal diurnal*.  $M_2$

tides in the Gulf of Alaska are larger than either  $O_1$  or  $K_1$ . However, this area (region 7) is still classified as predominantly *nodal diurnal*.

The situation is somewhat reversed in the western Atlantic Ocean, where semidiurnal tides dominate to such an extent that five of its seven listed regions are classified as *nodal semidiurnal*. Only the Gulf of Mexico (region B) is dominated by diurnal tides and is therefore *nodal diurnal*, while the Gulf of St. Lawrence (region D) is *nodal diurnal* for the same reasons as the Gulf of Alaska. The Northwest Atlantic Ocean has two areas which compete for the largest  $M_2$  tides in the world. These are Bay of Fundy (in region C) and Ungava Bay (in region F), where  $M_2$  amplitudes exceed 400 cm. It is therefore not surprising that these

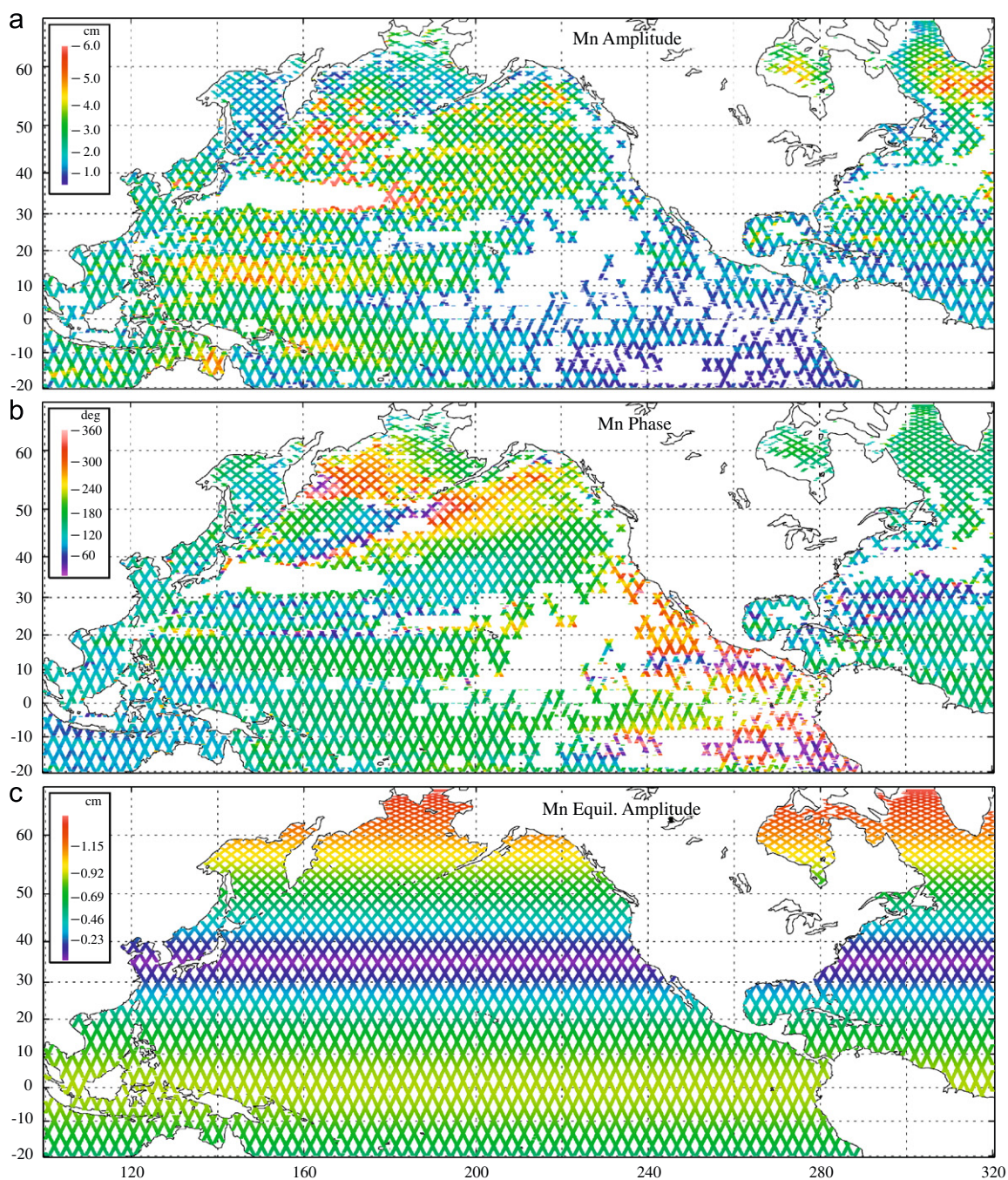


Fig. 6. Computed (a) amplitudes and (b) phases of  $M_n$  and (c) the equilibrium amplitude of  $M_n$ .

regions and the adjacent areas of Hudson Bay (region E) and northern Labrador Sea (region G) are predominantly *nodal semidiurnal*.

We also included in the analysis the 18.6-year constituent  $M_n$ . Its amplitudes (Fig. 6a) appear to be quite small, 1.5–3.5 cm over most of the area, and phase values (Fig. 6b) are clustered in the  $163 \pm 64^\circ$  range. Even after applying our edit criteria, which eliminated most of the areas of the Kuroshio, the Oyashio, the Gulfstream and their extensions, there are clear anomalies in  $M_n$  amplitudes and phases in the areas of enhanced decadal variability and eddy activity, for example, east of the Korean Peninsula (cf. Kang et al., 2005) and in the Oyashio Current and its extension area. Because of its small amplitude and the effects of decadal ocean variability, such as the Pacific Decadal Oscillation, the altimetry-derived  $M_n$  amplitudes and phase values likely arise from this variability and not the nodal tide.

The theoretical 18.6-year nodal tide is a function of latitude,  $0.9(3\sin^2\phi - 1)\cos N$  cm (e.g., Pugh, 1987), where  $N$  is the current longitude of the moon's ascending node. Its amplitude has a distinct zonal structure and is less than 1 cm almost everywhere (Fig. 6c). This is well below the level of natural ocean variability at this low frequency and the mean amplitude of the analysed  $M_n$ , which is about 2.5 cm (Table 1). Given the nonstationary nature of decadal ocean variability, our along track results for the  $M_n$  amplitudes from altimetry data covering less than one period do not contradict the results of Trupin and Wahr (1990), who were

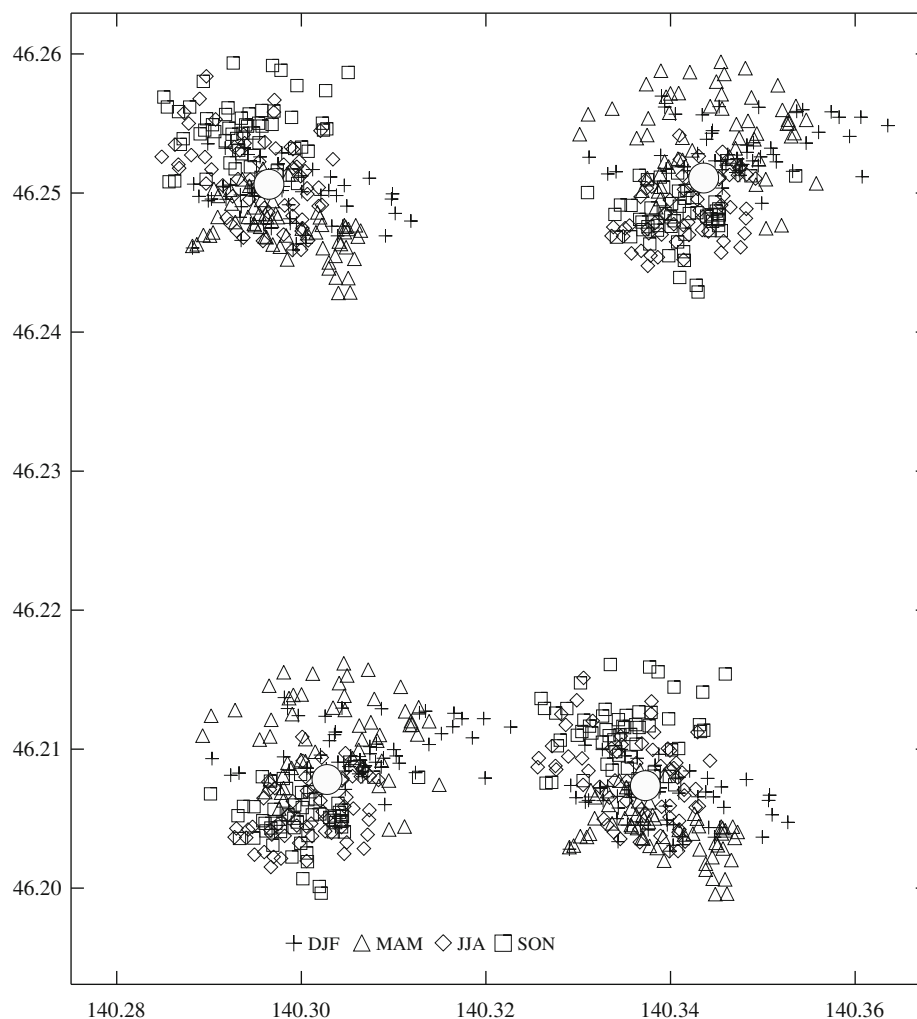
able to get a good fit to the theoretical 18.6-year spectral peak by “stacking” much longer time series of yearly averaged global data set of 260 coastal tide gauges.

## 5. Conclusions

In many parts of the ocean 16 years of collocated TPJ altimetry data are long enough for harmonic analyses to adequately separate the nodal satellites  $Q_{1n}$ ,  $O_{1n}$ ,  $K_{1n}$ ,  $M_{2n}$  and  $K_{2n}$  from their parent constituents  $Q_1$ ,  $O_1$ ,  $K_1$ ,  $M_2$  and  $K_2$ , even for the weaker satellites  $Q_{1n}$ ,  $M_{2n}$  and  $K_{2n}$ , with mean amplitudes between 1 and 2 cm. The utility of such analyses is evident from a relatively good agreement between the corresponding amplitude and phase patterns in Figs. 1–5 and from small off-diagonal elements in the error covariance matrices.

However, due to the relatively small amplitudes of these satellites (Table 1) and the likely influence from low-frequency aliased signals, the observed amplitude ratios between the nodal satellites and their parent constituents tend to exceed the values predicted from equilibrium tidal theory. This amplitude excess tends to be minimal where parent amplitudes are relatively large, for example, in the Sea of Okhotsk for diurnal and in the Yellow and East China Seas for semidiurnal constituents.

Examination of diurnal and semidiurnal nodal amplitudes in select coastal areas and marginal seas around the Pacific and the



**Fig. A1.** Four data clusters near a crossover of Jason-1 passes 60 and 203 after along-track interpolation to 1-s time intervals. Four symbols show data locations during the months of winter (DJF), spring (MAM), summer (JJA) and fall (SON). Solid white circles mark mean “collocated” coordinates assigned to each data cluster.

western Atlantic Oceans allowed the assignment of a *nodal character* to these regions, which were classified as *nodal diurnal*, *nodal semidiurnal*, or *nodal mixed*, based on the nodal amplitudes in each band. While the areas with predominant diurnal tides are all *nodal diurnal*, due to their relatively large (0.19 and 0.14) nodal amplitude ratios, the small nodal ratio of 0.037 for  $M_{2n}$  resulted in some regions with strong  $M_2$  tides to be classified as *nodal diurnal* or *nodal mixed*. The amplitude ratio between  $K_{2n}$  and  $K_2$  is 0.30, making the  $K_{2n}$  amplitudes sometimes comparable to those of  $M_{2n}$ . However, this effect was not sufficient to make all of the areas with dominant  $M_2$  to be dominant *nodal semidiurnal*.

The observed amplitudes of the 18.6-year constituent  $M_n$  are relatively small (Fig. 6). However, they significantly exceed theoretical nodal tide amplitudes at these latitudes. The analysed  $M_n$  are therefore of non-tidal origin, likely due to broad-band ocean variability with non-stationary phases at this (18.6 year) $^{-1}$  frequency.

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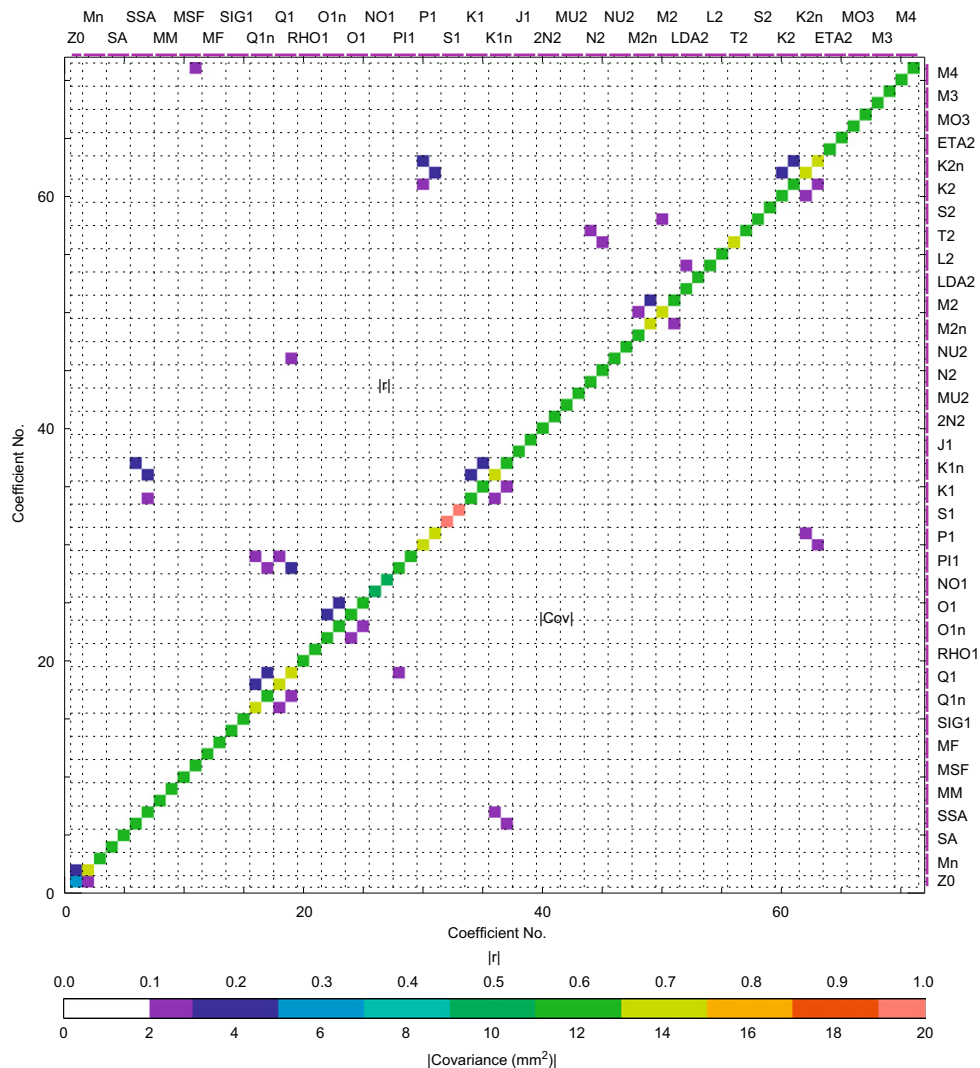
by the Korea Ocean Research and Development Institute (KORDI research program PG47100) and by Fisheries and Oceans Canada.

## Appendix A. Collocation of along-track data

For each altimeter cycle and pass in the selected area, the along-track data were first filtered for small-scale noise by applying twice an along-track 3-point median filter. The filtered data were then interpolated to exact 1-s time intervals (ground track separation of about 5.8 km), counting from the equator crossing for ascending and to equator crossing for descending passes. Because of seasonal and longer-term variations in satellite orbits, these 1-s samples are actually spread over a small area, forming patches with nominal (rms) radii of the order of 1 km at mid-latitudes (Fig. A1). For the purposes of this study, we consider these patches to be sufficiently small and thus “collocated” in space.

## Appendix B. Constituent error-covariance

CFCH discussed the use of error covariance matrices, computed from a least-squares solution of the overdetermined problem, for



**Fig. B1.** An example of the error covariance matrix  $c_{ij}$  at an along-track location south of the Korean Peninsula (34.4N, 128.7E). As this matrix is symmetrical,  $|c_{ij}|$  values are shown below and on the diagonal and correlation coefficients  $|r_{ij}|$  above the diagonal.

**Table B1**

Aliased constituent pairs which require more than 3400 days to separate.

	$Q_1 : Q_{1n}$	$O_1 : O_{1n}$	$P_1 : K_2$	$P_1 : K_{2n}$	$K_1 : S_{sa}$	$K_1 : K_{1n}$	$K_{1n} : S_{sa}$	$M_2 : M_{2n}$	$K_2 : K_{2n}$
$R$	0.88	0.88	1.75	0.88	1.75	0.88	0.88	0.88	0.88
$T_{min}$	6798	6798	3404	6820	3404	6798	6798	6798	6796

$R$  is the Rayleigh number and  $T_{min}$  (in days) is a minimum period required to resolve the two constituents, based on TPJ cycle period of 9.916 days and time series duration of 5970 days.

additional guidance on selection of tidal constituents and for assessment of aliasing and validity of local solutions. Fig. B1 shows an example of the error covariance matrix  $c_{ij} = \text{cov}(x_i, x_j)$ , where  $x_i$  ( $i = 1, \dots, 2K - 1$ ) is the solution vector of sine and cosine coefficients. For  $K = 36$  tidal constituents listed in Table 1, this matrix has  $71 \times 71$  elements.

As this matrix is symmetrical,  $|c_{ij}|$  values are shown in Fig. B1 below and on the diagonal, while correlation coefficients  $|r_{ij}| = |c_{ij}|/(c_{ii}c_{jj})^{1/2}$  are shown above the diagonal. At this particular location (34.43N, 128.74E), the number of valid observations is  $M = 538$  (out of possible 624) and the maximum error variance values on the diagonal are about  $17 \text{ mm}^2$  for the two coefficients of  $S_1$ . In general, the error covariances are quite small for this 16.3-year long data set, compared to their values in CFCH, where  $M < 250$ .

For a tidal constituent  $k$ , an estimate of its amplitude error was defined in CFCH from their Eq. (6) as  $\sigma_k = (c_{2k-2,2k-2} + c_{2k-1,2k-1})^{1/2}$  for  $k = 2, \dots, K$  and  $\sigma_1 = c_{1,1}^{1/2}$  for  $Z_0$ . This was redefined here as  $\sigma_k = (\sum_{i=1,2K-1} (c_{i,2k-2}^2 + c_{i,2k-1}^2)^{1/2})^{1/2}$ , thus also including the off-diagonal elements, though over most of the area the off-diagonal terms are relatively small, with their combined contribution to  $\sigma_k$  usually less than 20%. Only a few off-diagonal  $|c_{ij}|$  in Fig. B1 approach  $3\text{--}4 \text{ mm}^2$ , for example,  $P_1\text{--}K_{2n}$  and  $S_{sa}\text{--}K_{1n}$ .

The  $|r_{ij}|$  values in this figure are somewhat larger between the 18.6-year satellites and their parent constituents, as well as for  $S_{sa}\text{--}K_{1n}$ ,  $P_1\text{--}K_{2n}$  and  $Z_0\text{--}M_n$  pairs. This is as expected, as for these pairs the Rayleigh criterion is  $R < 1$ , for a TPJ time series with a duration of approximately 5970 days, while the minimum period  $T_{min}$  required to resolve all the nodal constituents is about 6800 days (Table B1).

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